

Gravitational waves

The background of the slide is a 3D visualization of a gravitational well. It features a grid of blue and green lines that curve and ripple, representing the curvature of spacetime. In the center-right of the image, two bright blue spheres, representing black holes, are shown in a binary system, with the grid curving around them.

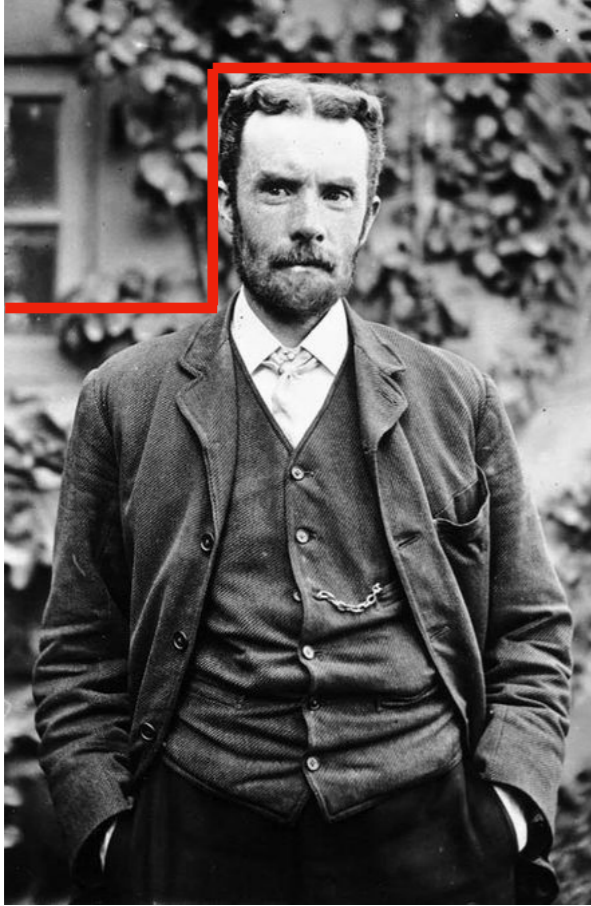
A new window on the Universe

Adrien Kuntz

L'Agape
July 2019

Part I : a short history

A bit of history

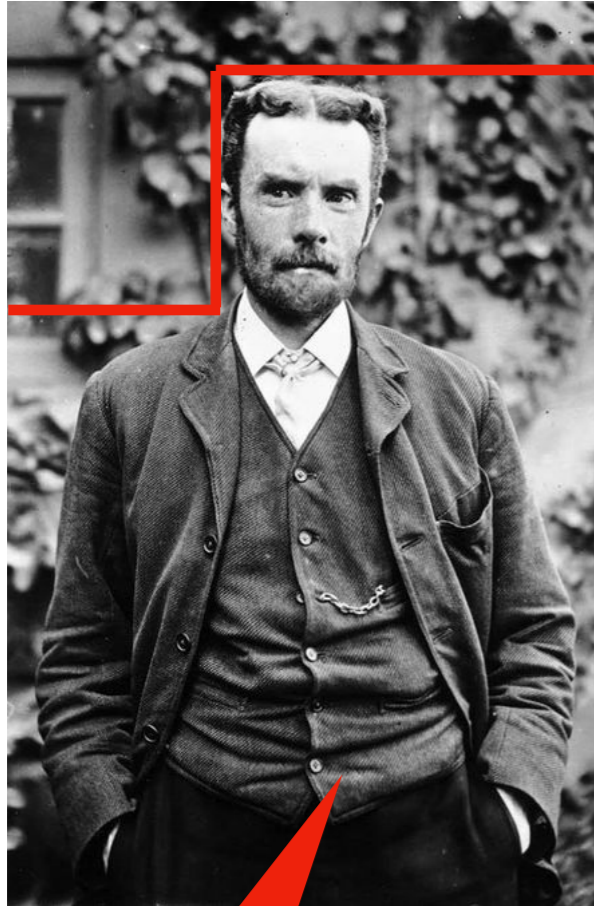


$$v^2 \nabla^2 e = \frac{\partial^2 e}{\partial t^2},$$



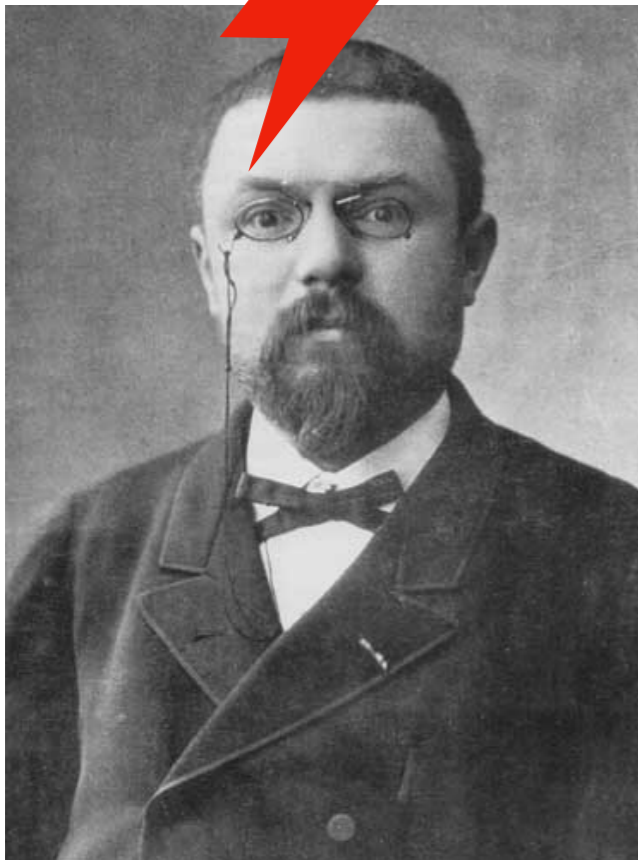
Quand nous parlerons donc de la position ou de la vitesse du corps attirant, il s'agira de cette position ou de cette vitesse à l'instant où l'*onde gravifique* est partie de ce corps; quand nous parlerons de la position ou de la vitesse du corps attiré, il s'agira de cette position ou de cette vitesse à l'instant où ce corps attiré a été atteint par l'onde gravifique émanée de l'autre corps; il est clair que le premier instant est antérieur au second.

A bit of history



$$v^2 \nabla^2 \mathbf{e} = \frac{\partial^2 \mathbf{e}}{\partial t^2},$$

Mathematics is an experimental science, and definitions do not come first, but later on. They make themselves, when the nature of the subject has developed itself



Quand nous parlerons donc de la position ou de la vitesse du corps attirant, il s'agira de cette position ou de cette vitesse à l'instant où l'onde gravifique est partie de ce corps; quand nous parlerons de la position ou de la vitesse du corps attiré, il s'agira de cette position ou de cette vitesse à l'instant où ce corps attiré a été atteint par l'onde gravifique émanée de l'autre corps; il est clair que le premier instant est antérieur au second.

A bit of history



$$h = \frac{2G}{c^4} \frac{1}{r} \frac{\partial^2 Q}{\partial t^2}$$

Transverse-transverse h_{22}, h_{33}, h_{23} .
Longitudinal-transverse $h_{12}, h_{13}, h_{24}, h_{34}$.
Longitudinal-longitudinal h_{11}, h_{14}, h_{44} .



They are not objective, and (like absolute velocity) are not detectable by any conceivable experiment. They are merely sinuosities in the co-ordinate-system, and the only speed of propagation relevant to them is "the speed of thought."

A bit of history



$$h = \frac{2G}{c^4} \frac{1}{r} \frac{\partial^2 Q}{\partial t^2}$$

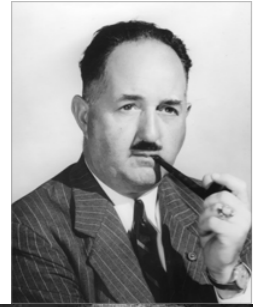
Transverse-transverse	h_{22}, h_{33}, h_{23} .
Longitudinal-transverse	$h_{12}, h_{13}, h_{24}, h_{34}$.
Longitudinal-longitudinal	h_{11}, h_{14}, h_{44} .

Together with a young collaborator, I arrived at the interesting result that gravitational waves do not exist, though they had been assumed a certainty to the first approximation.

They are not objective, and (like absolute velocity) are not detectable by any conceivable experiment. They are merely sinuities in the co-ordinate-system, and the only speed of propagation relevant to them is "the speed of thought."



To cut a long story short...



Robertson (referee) to Einstein and Rosen : your conclusion an artifact of your coordinate system



Einstein : I will never publish in Physical Review again



Infeld : Look, Einstein, Robertson was right

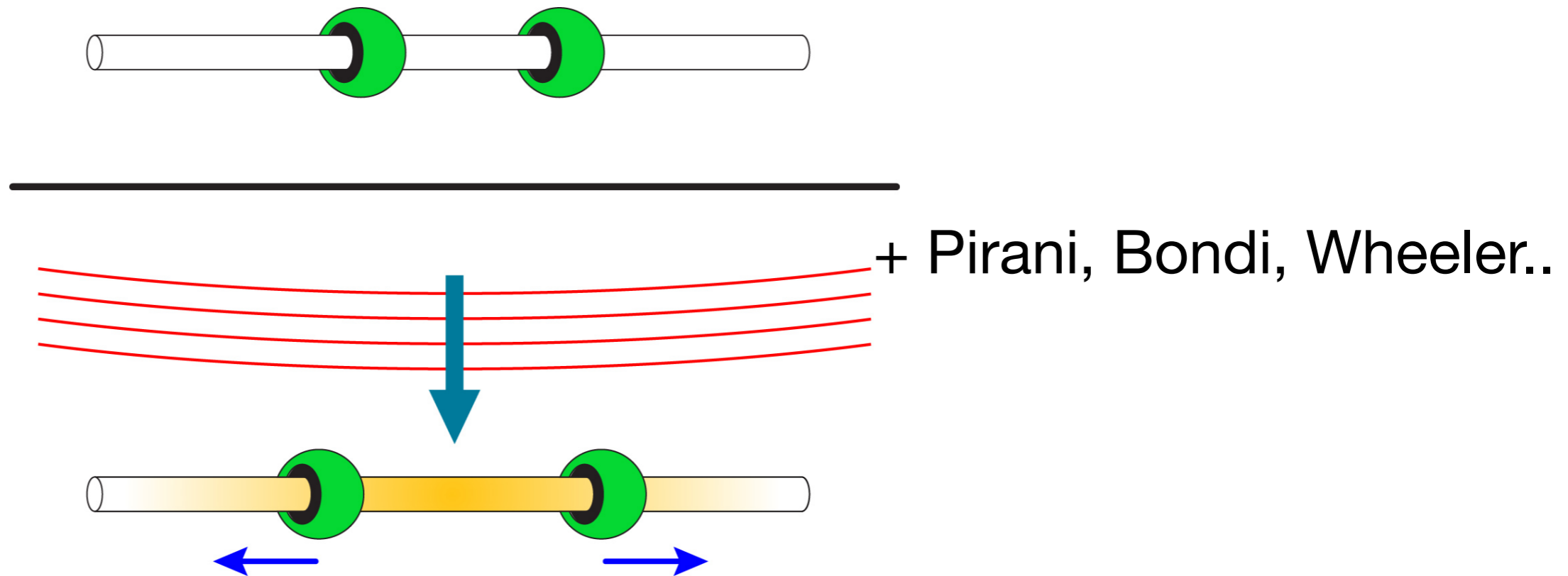


Einstein to Rosen (1937) : OK, so GW really exist. Let's just change the conclusion of our article.



Rosen : I still don't believe it and I prefer to publish my own version

Later on (1955)



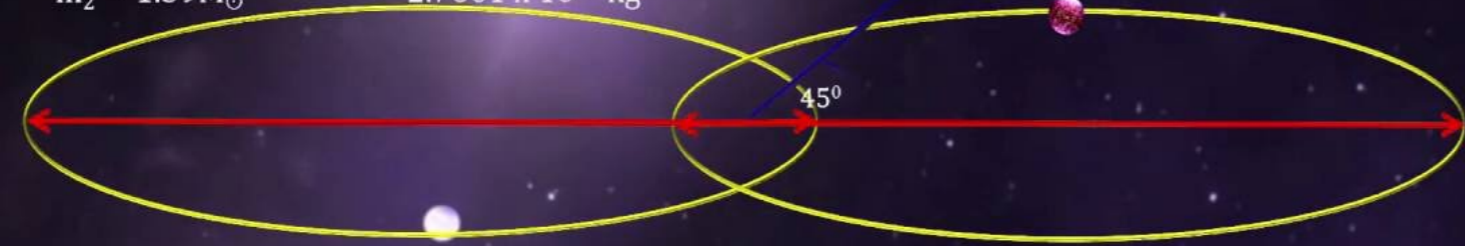
Rosen still did not believe until the 70's...

(And even today I heard strange statements from a physics teacher in Marseille, whose name I will keep secret...)

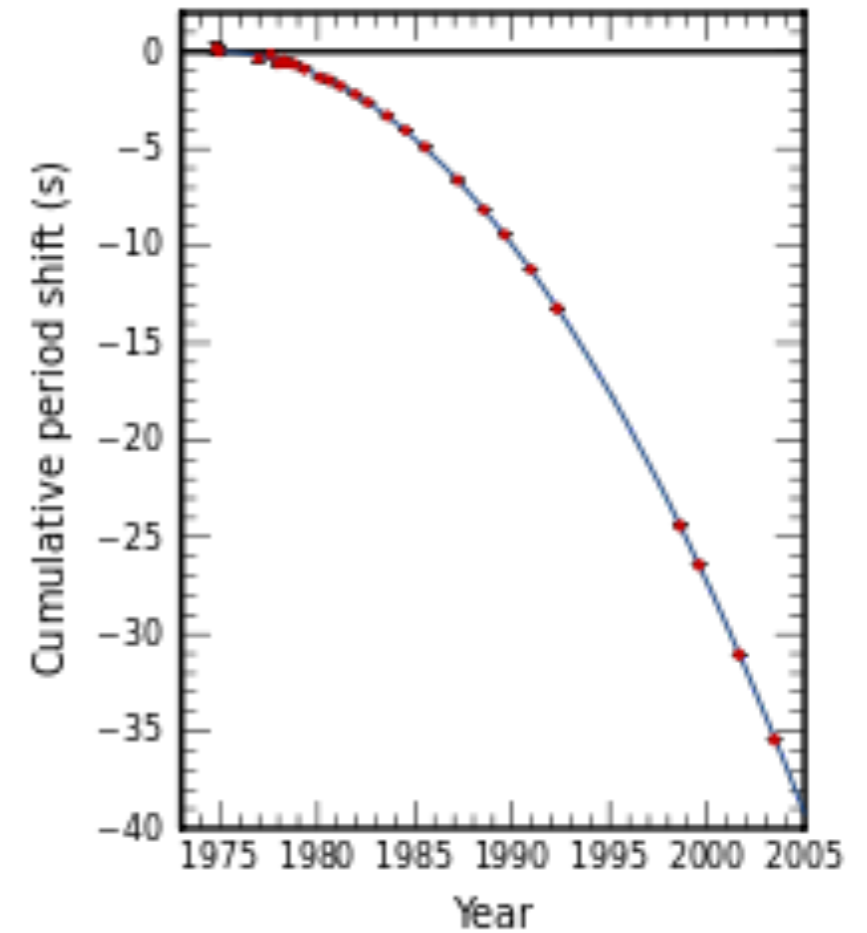
The Hulse-Taylor pulsar

PSR B1913+16

T = orbital period = 7.751939106 hr
 a = semi-major axis = 1.95×10^6 m
 e = eccentricity = 0.617131
 $m_1 = 1.44M_{\odot} = 2.8676 \times 10^{30}$ kg
 $m_2 = 1.39M_{\odot} = 2.7661 \times 10^{30}$ kg



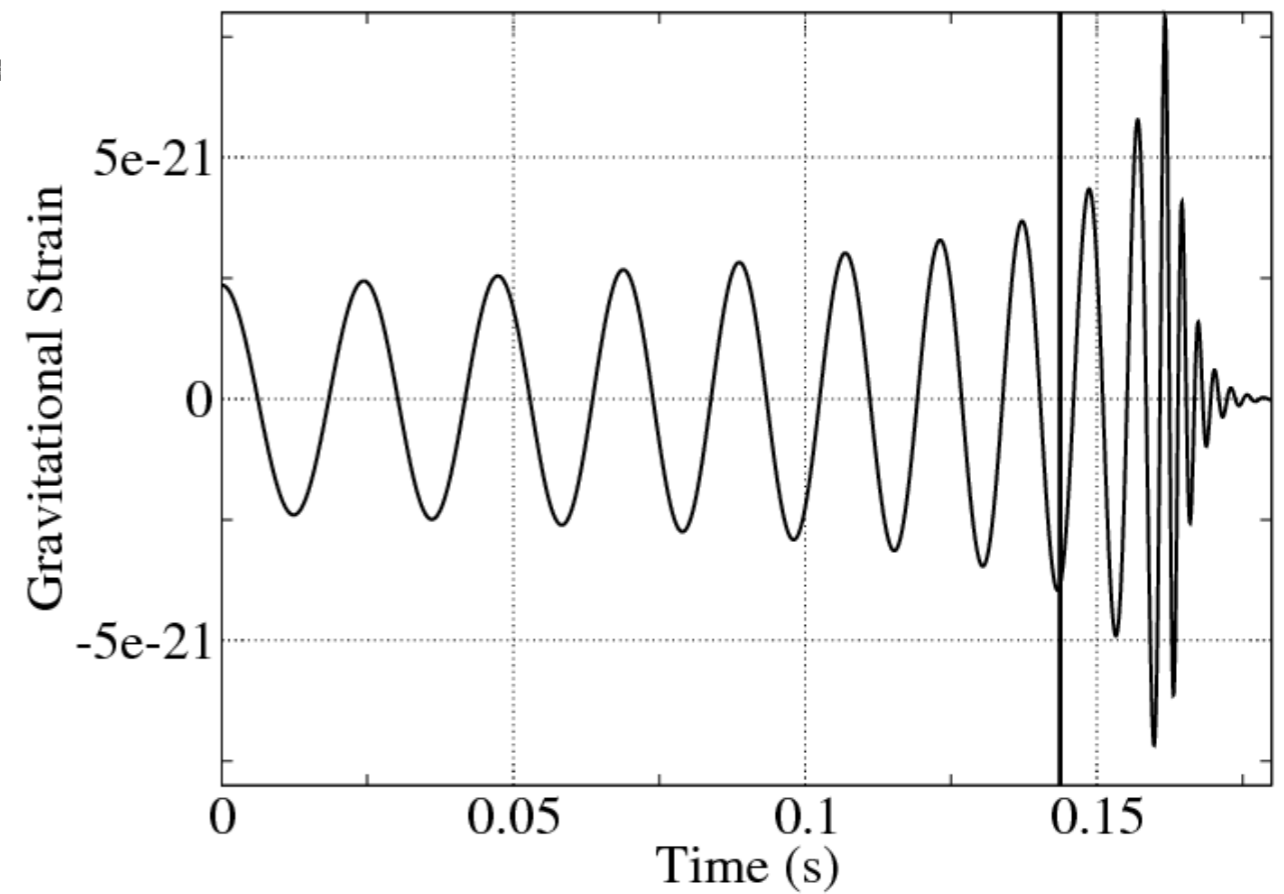
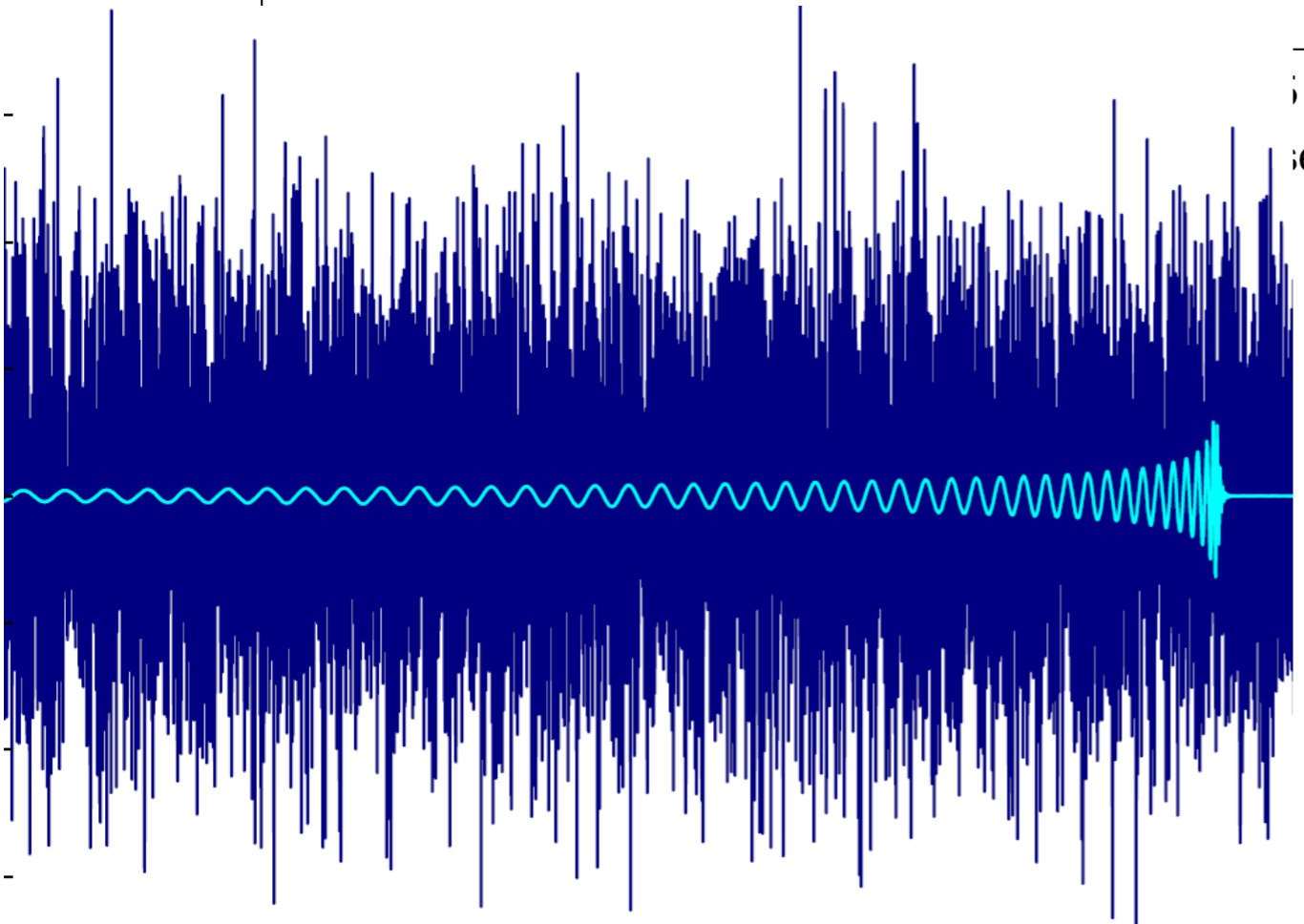
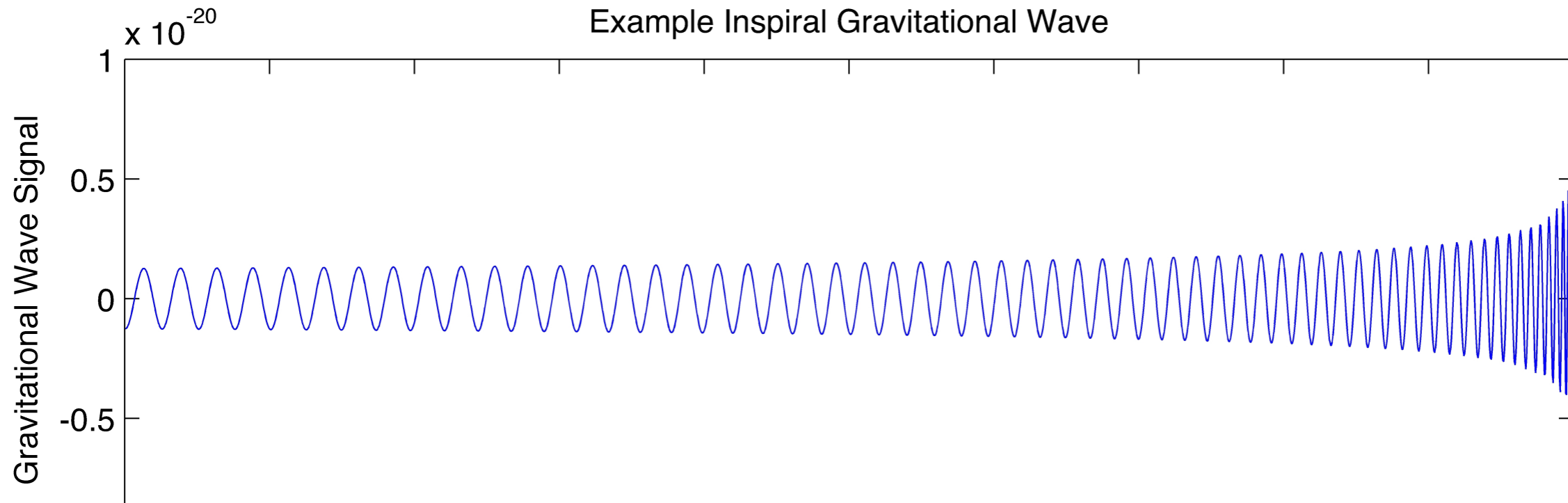
Periastron = 0.746×10^6 km
Apastron = 3.153×10^6 km
Inclination = 45°



Consistent with GR at the 0.2% level !

Today

Example Inspiral Gravitational Wave



Part II : sources of GW

Landscape

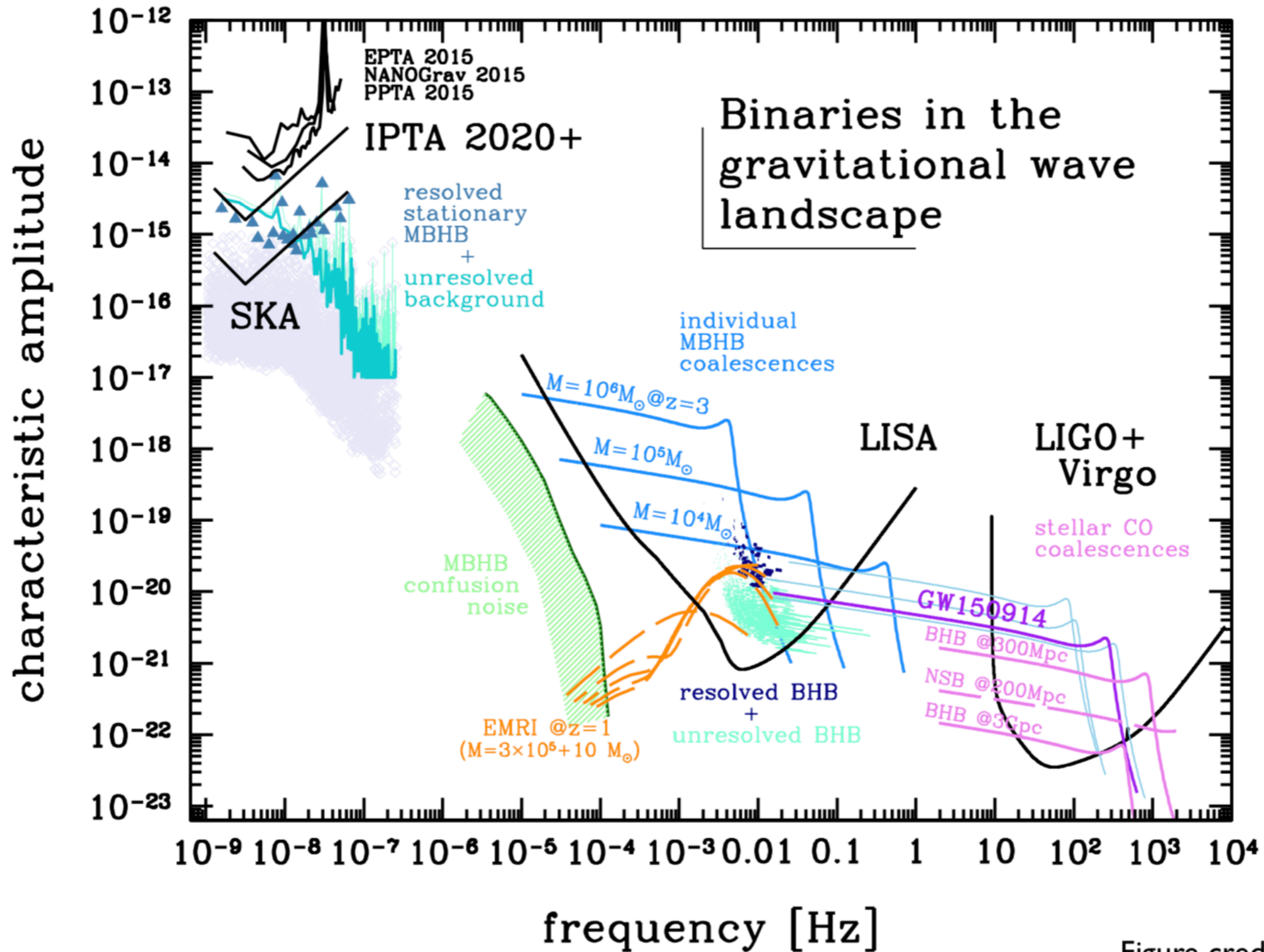
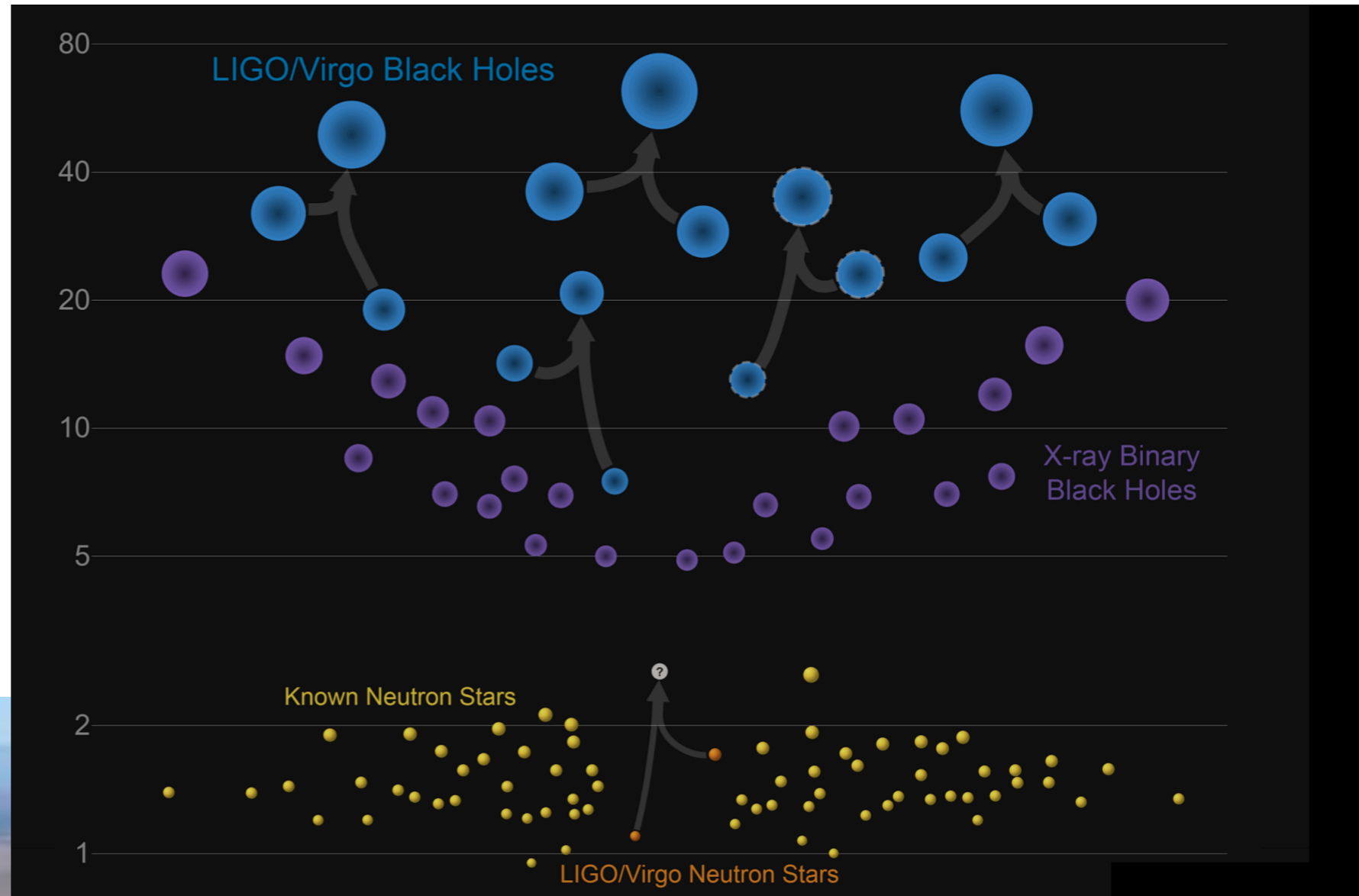


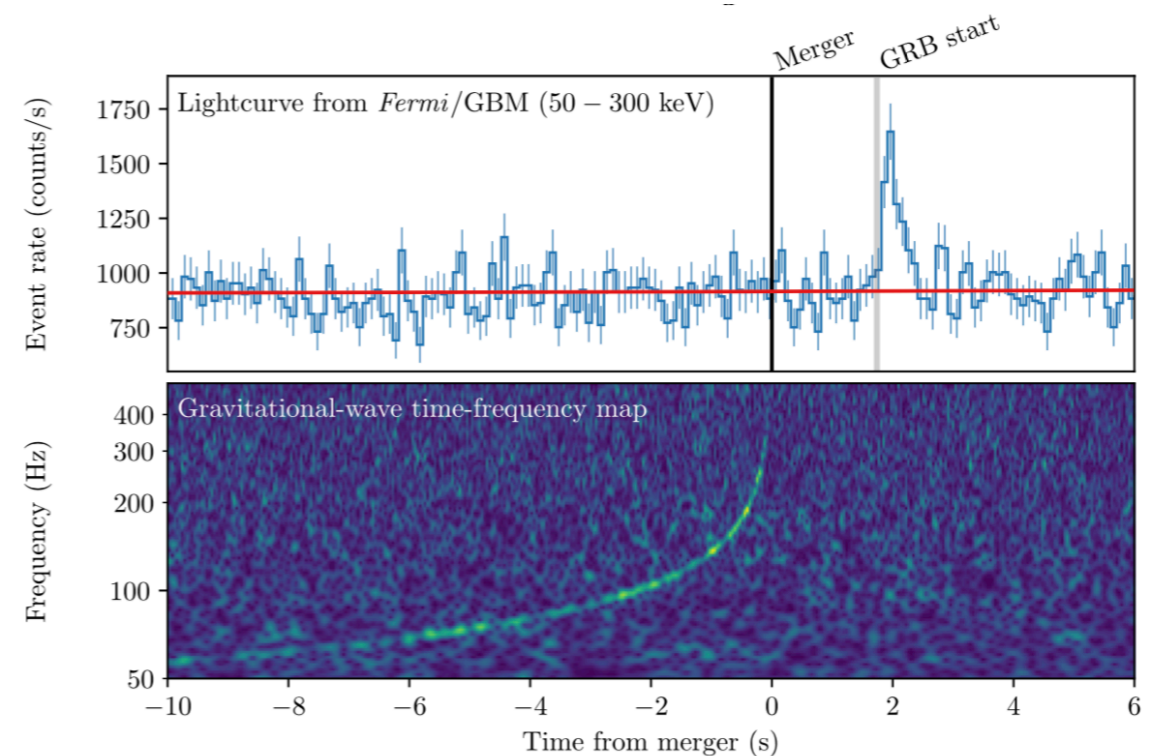
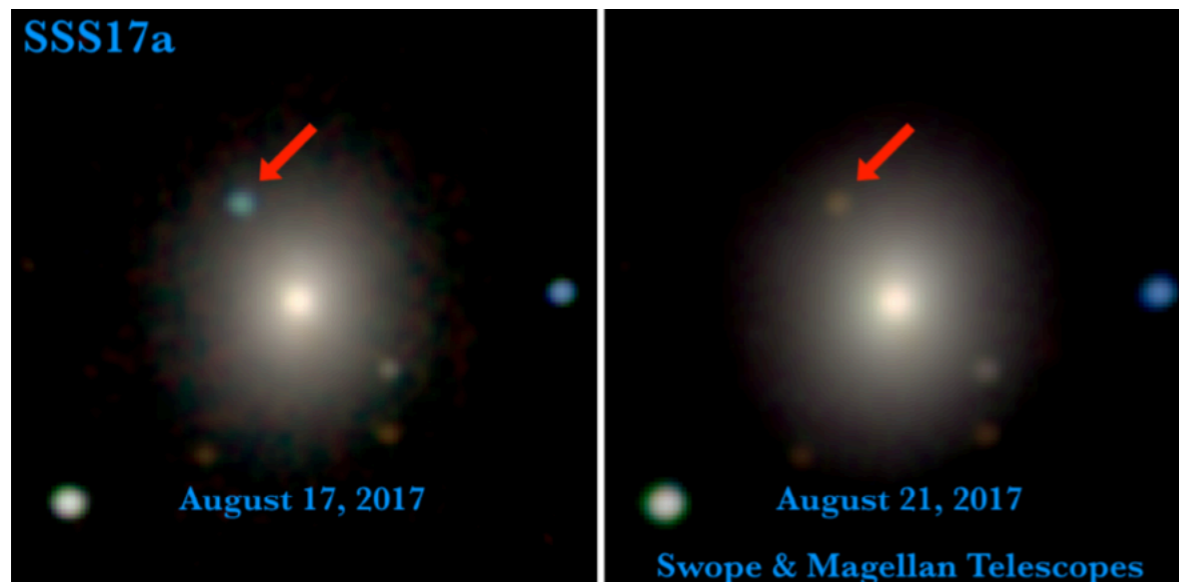
Figure credit: Alberto Sesana

LIGO/VIRGO



LIGO sources

- Most sources at $z \sim 0.1$ (less than 1 Gpc) : close universe
- BH-BH : spin and mass distribution, formation (star or primordial ?), no EM counterpart ?
- BH-NS : expected but not found yet
- NS-NS : 1 detection so far. EM counterpart : EOS of NS, speed of gravity, kilonova !

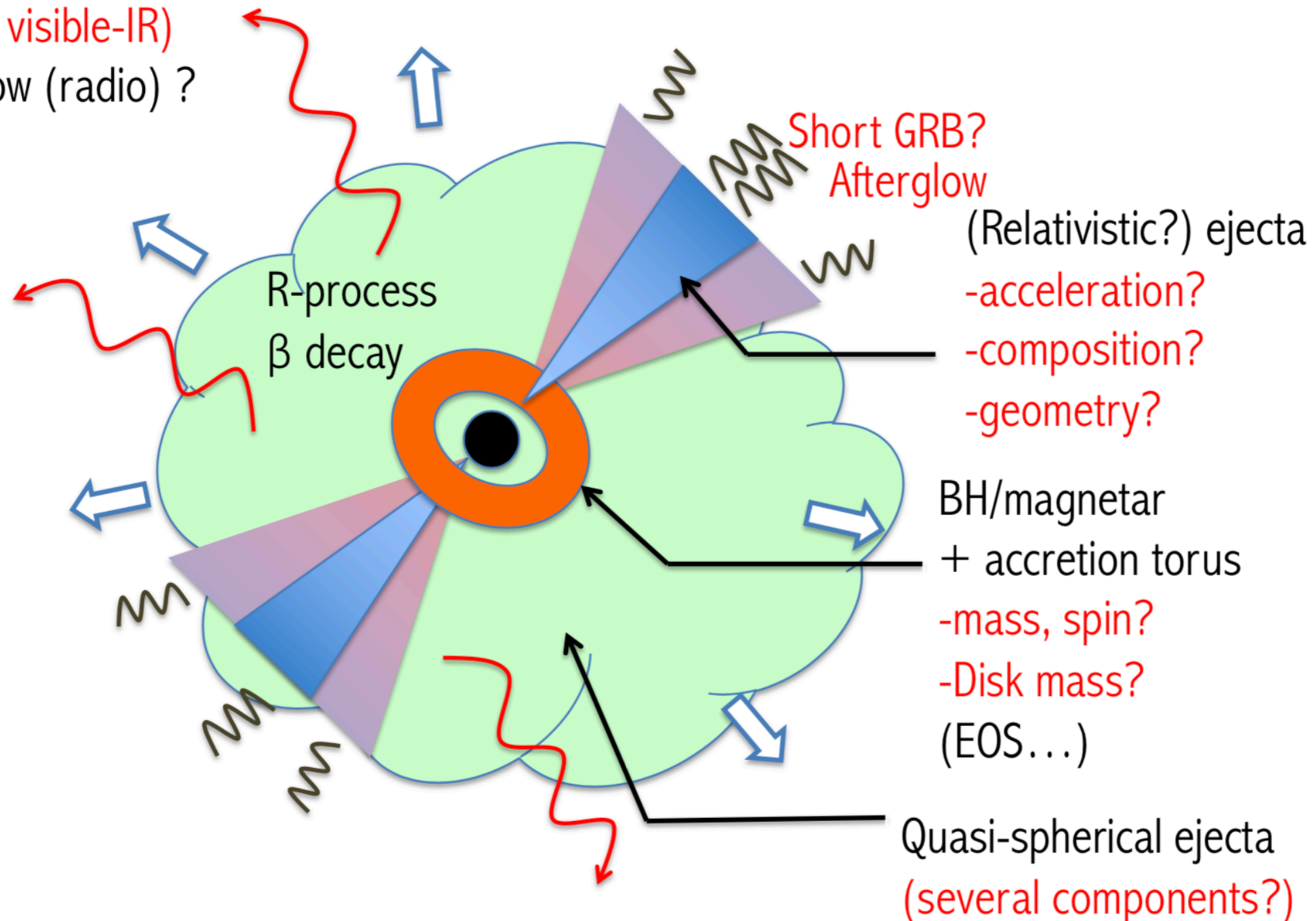


LIGO sources

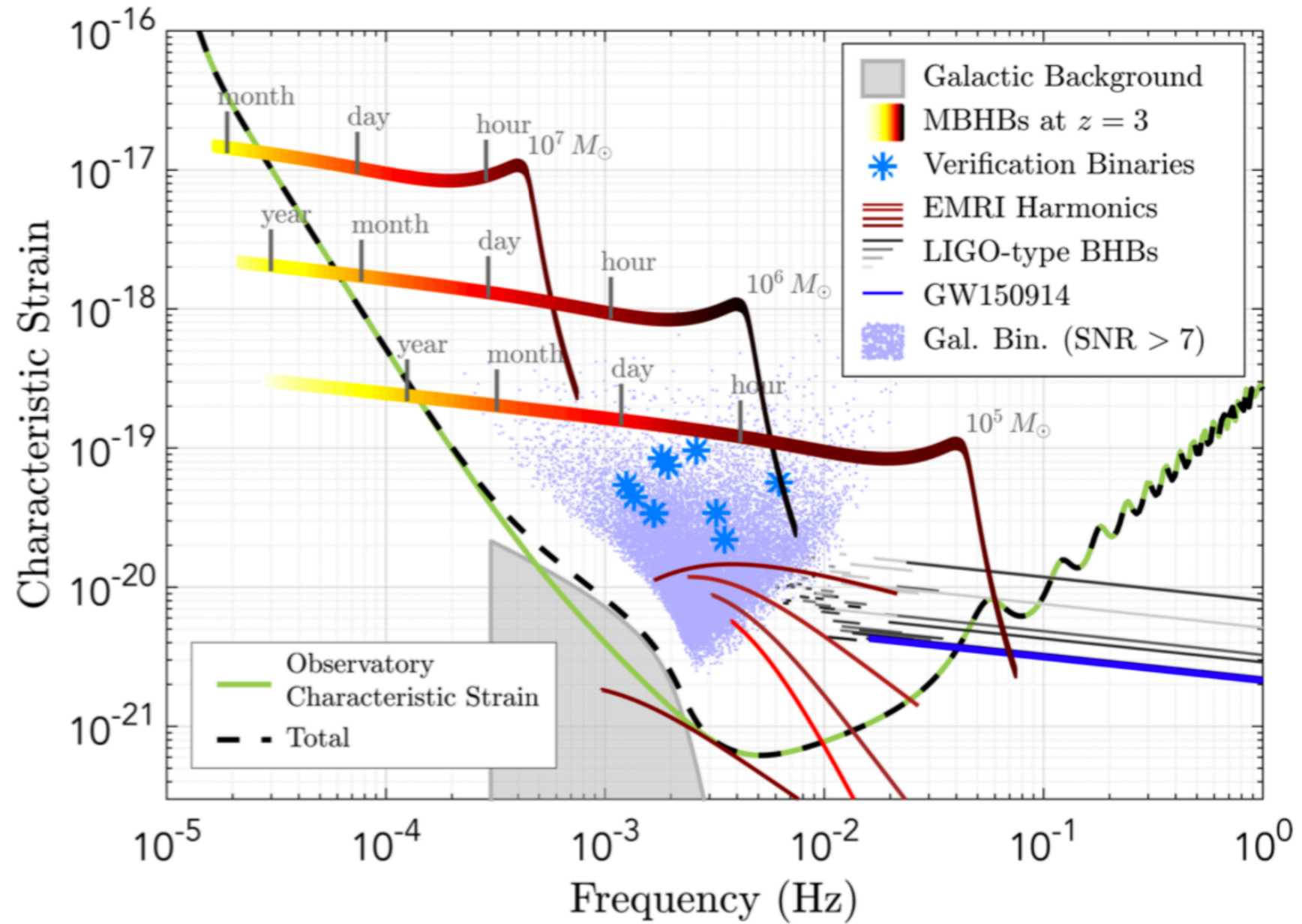
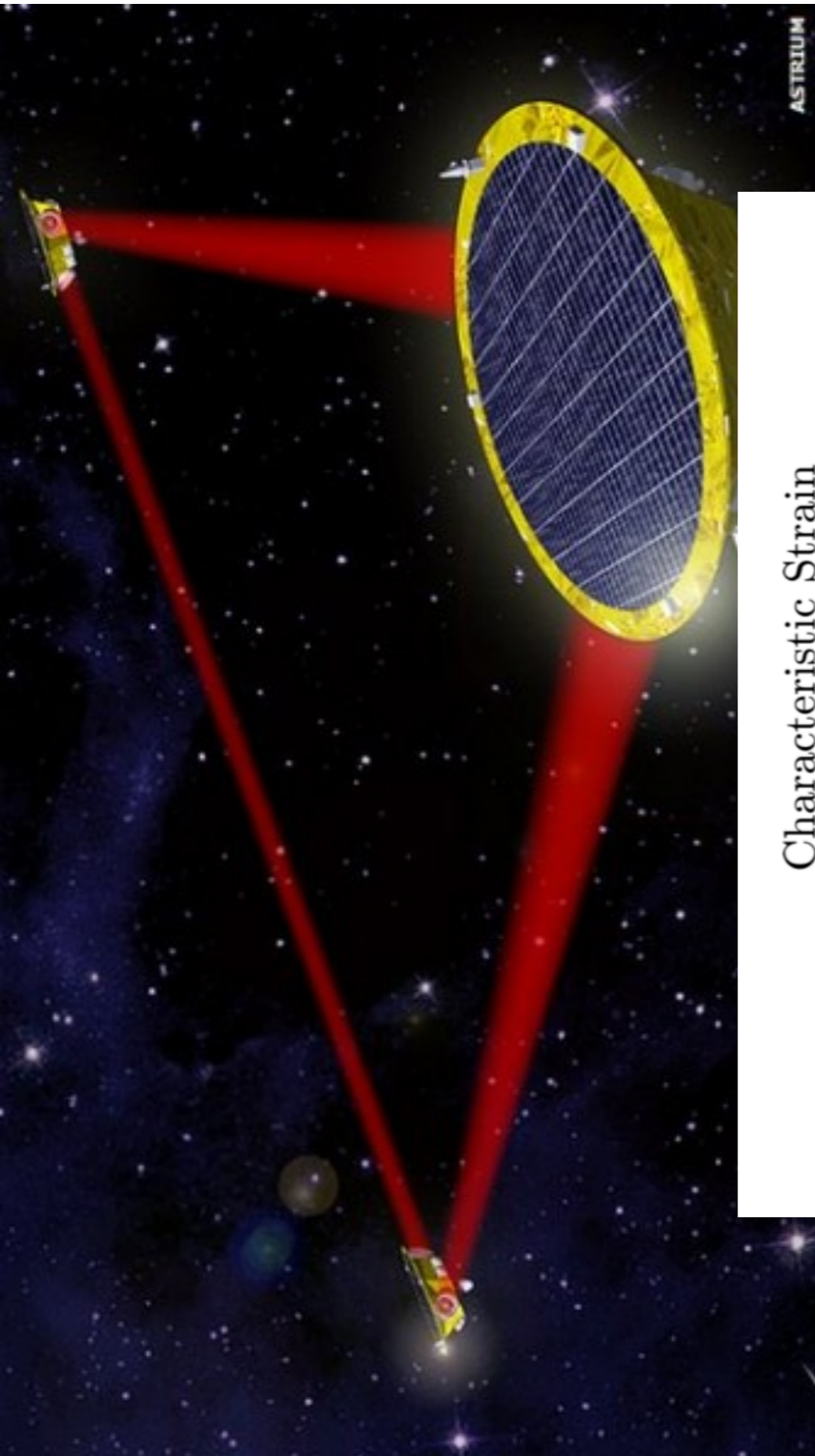
Remnant of a NS+NS merger

(Adapted from Frédéric Daigne)

Radioactively powered emission
(kilonova: visible-IR)
+ afterglow (radio) ?



LISA



(Credit : Maria Volonteri)

LISA sources

- Binaries in our Galaxy (background noise + 25000 resolved), before they enter the LIGO band !
- Supermassive BH merger, up to $z \sim 15 \Rightarrow$ population models, observation of the formation of a quasar in real time, H_0 measurement and cosmology,... Few events per year
- Extreme mass ratio inspiral ($60 M_{\odot}$ VS $10^5 M_{\odot}$) up to $z=4 \Rightarrow$ Very accurate measurement of the spin, eccentricity, inclination and test of gravity. Few events per year
- Stochastic GW background (inflation, cosmic strings...)
- Core collapse supernovae
- Exotic and unmodeled sources !

Some questions LISA will try to answer

How many galaxies host MBHs

→ when, where, how they form

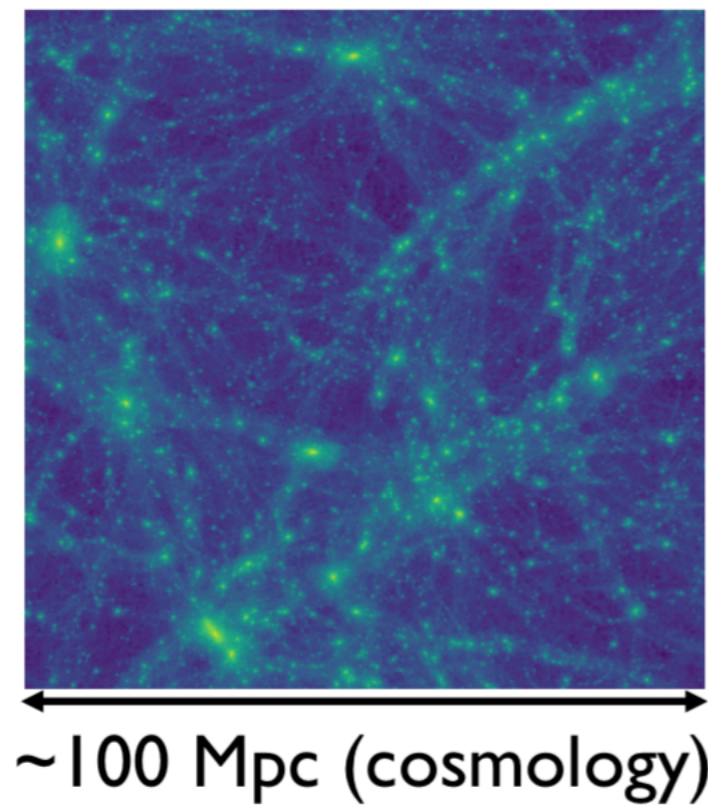
How long it takes for MBHs to merge in halo/
galaxy merger

→ dynamics of MBHs in mergers

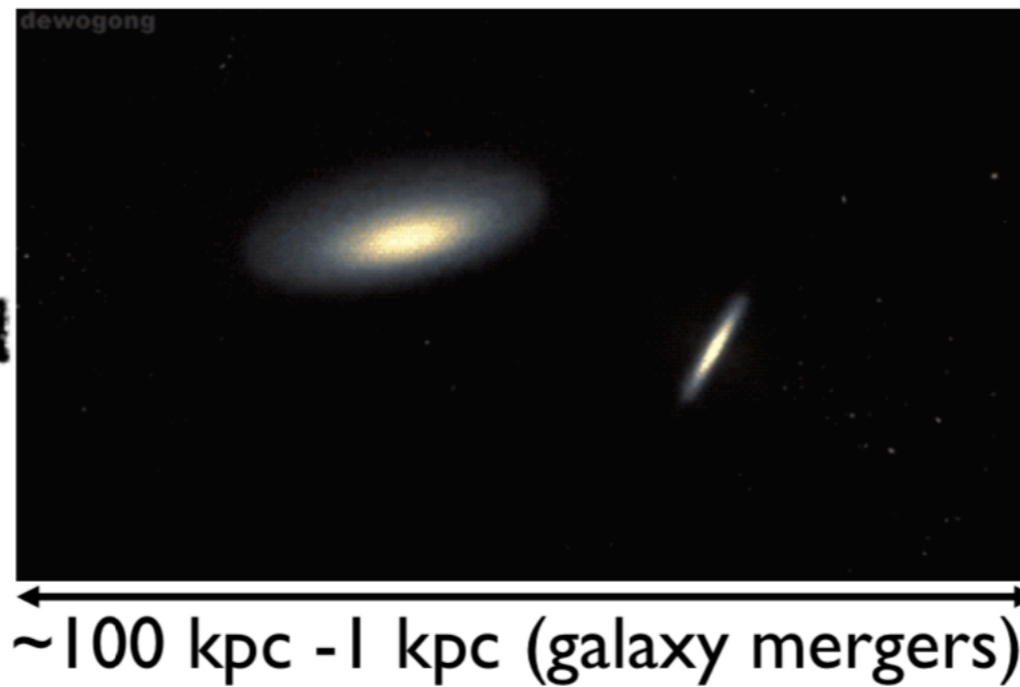
How MBHs grow in mass over time

→ accretion vs MBH-MBH mergers

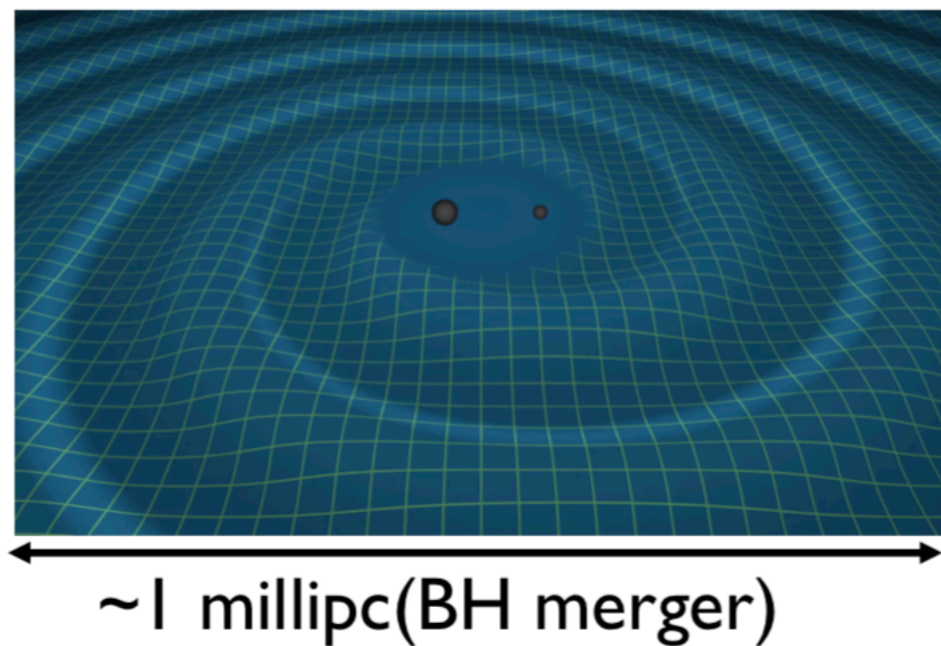
Some questions LISA will try to answer



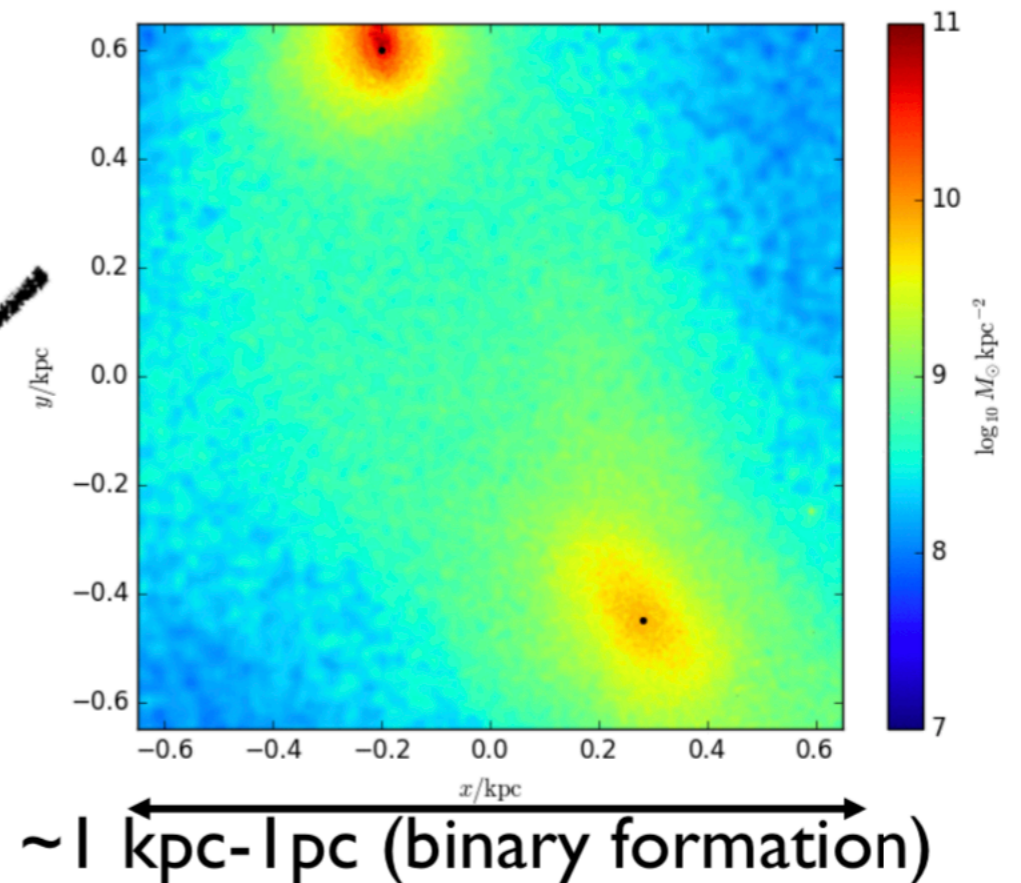
Gravity



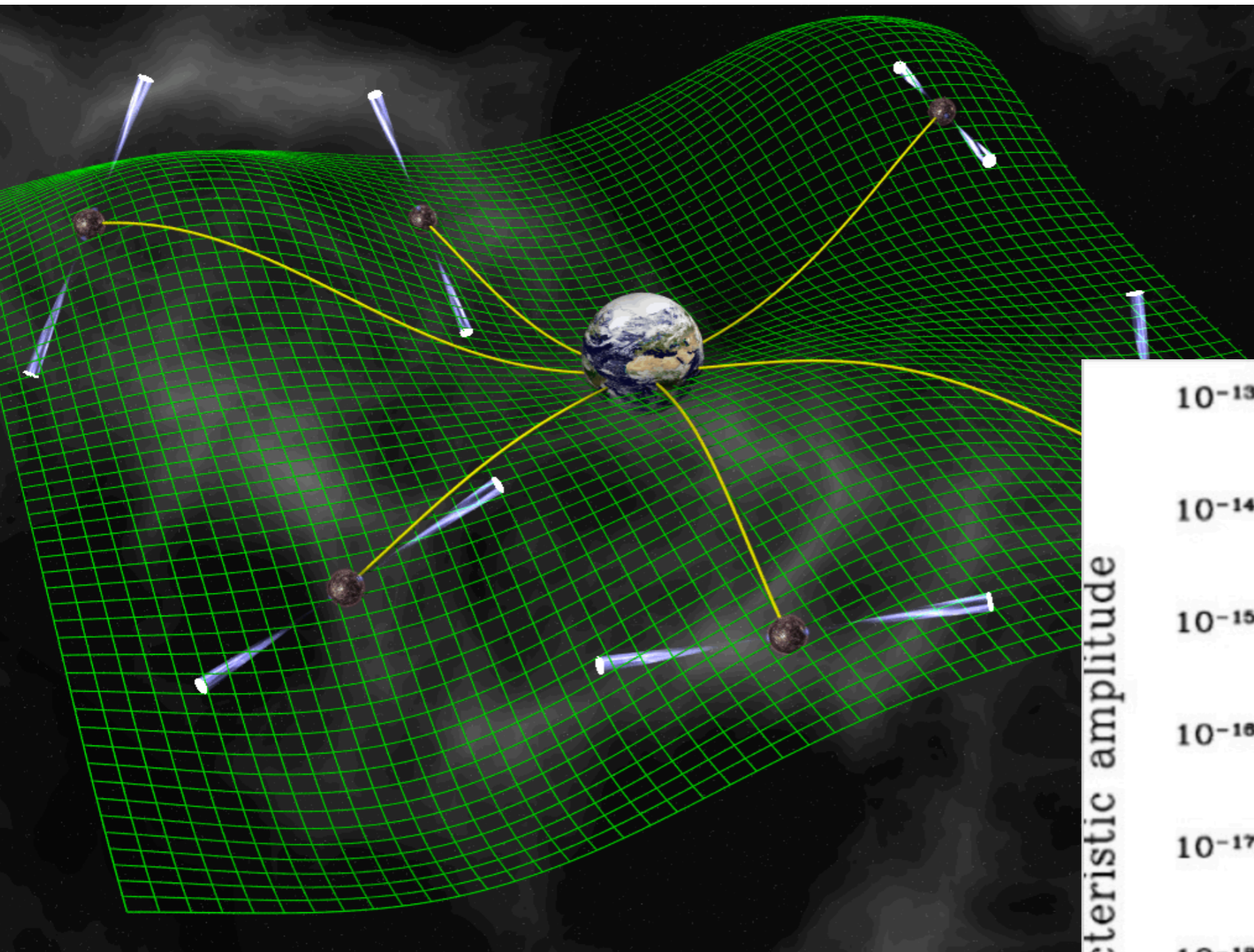
Dynamical Friction



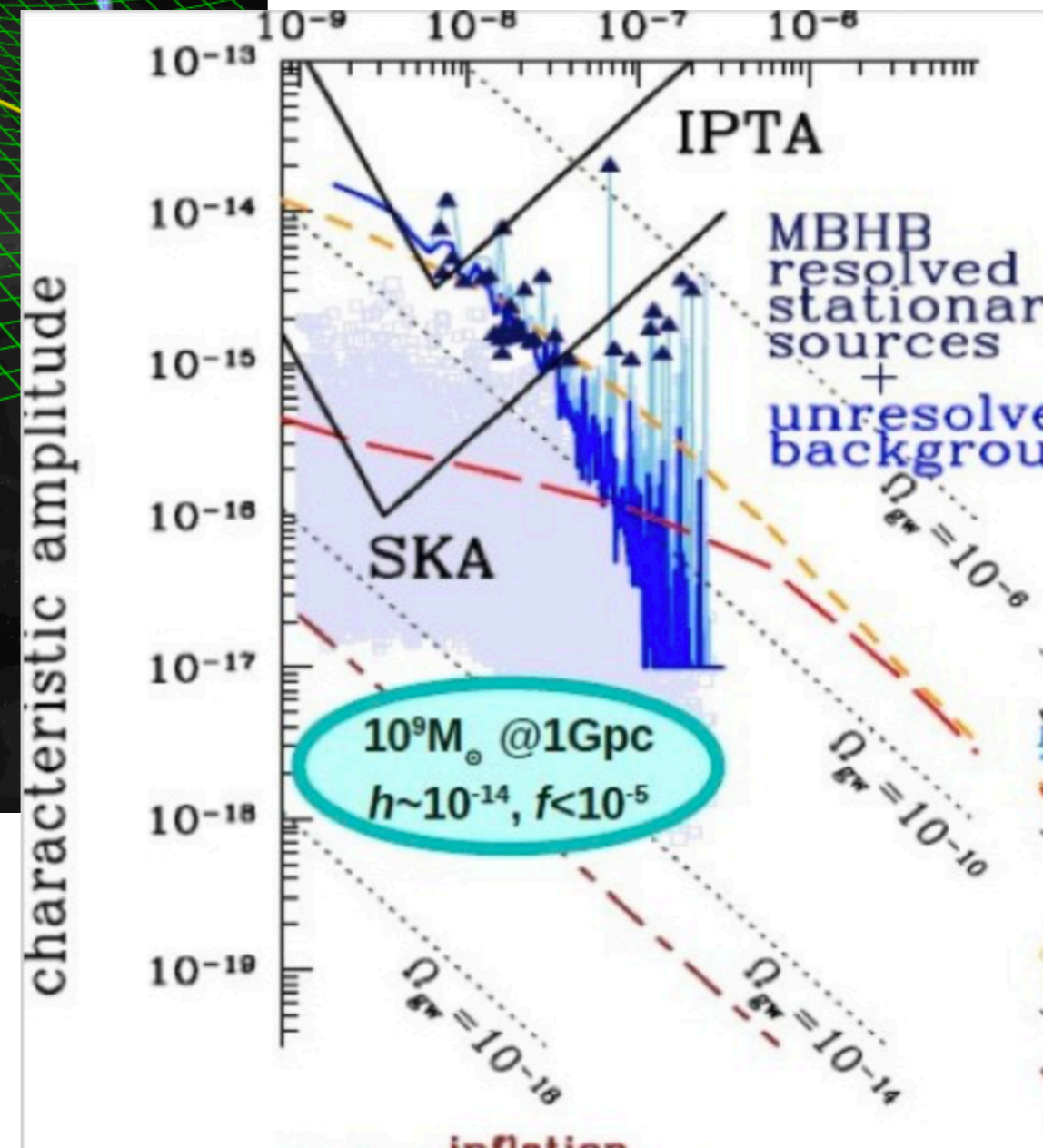
Gas torques?
Stellar scattering?
Last pc problem



Pulsar Timing Array

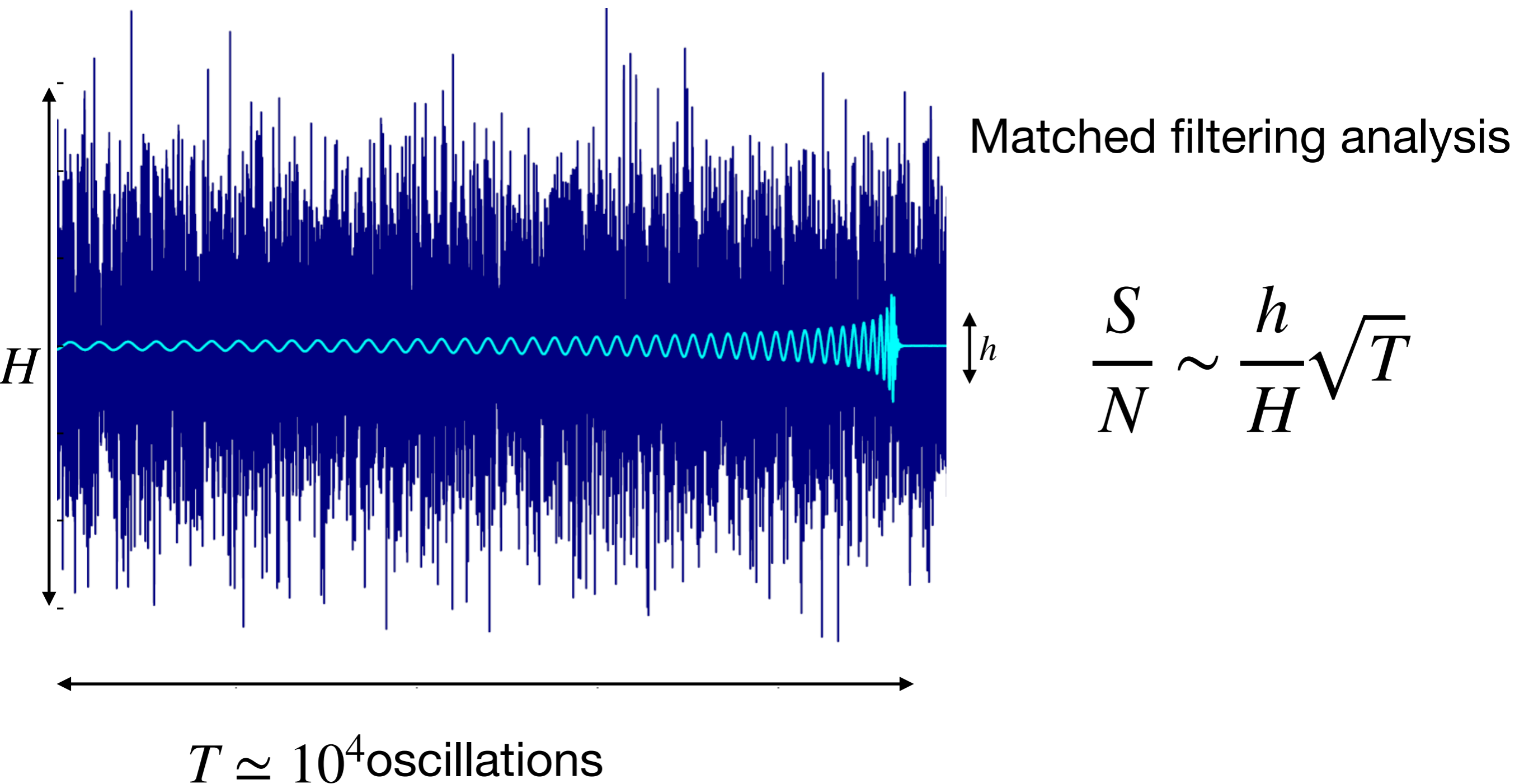


- Supermassive BH binaries
- Inflation
- Cosmic strings

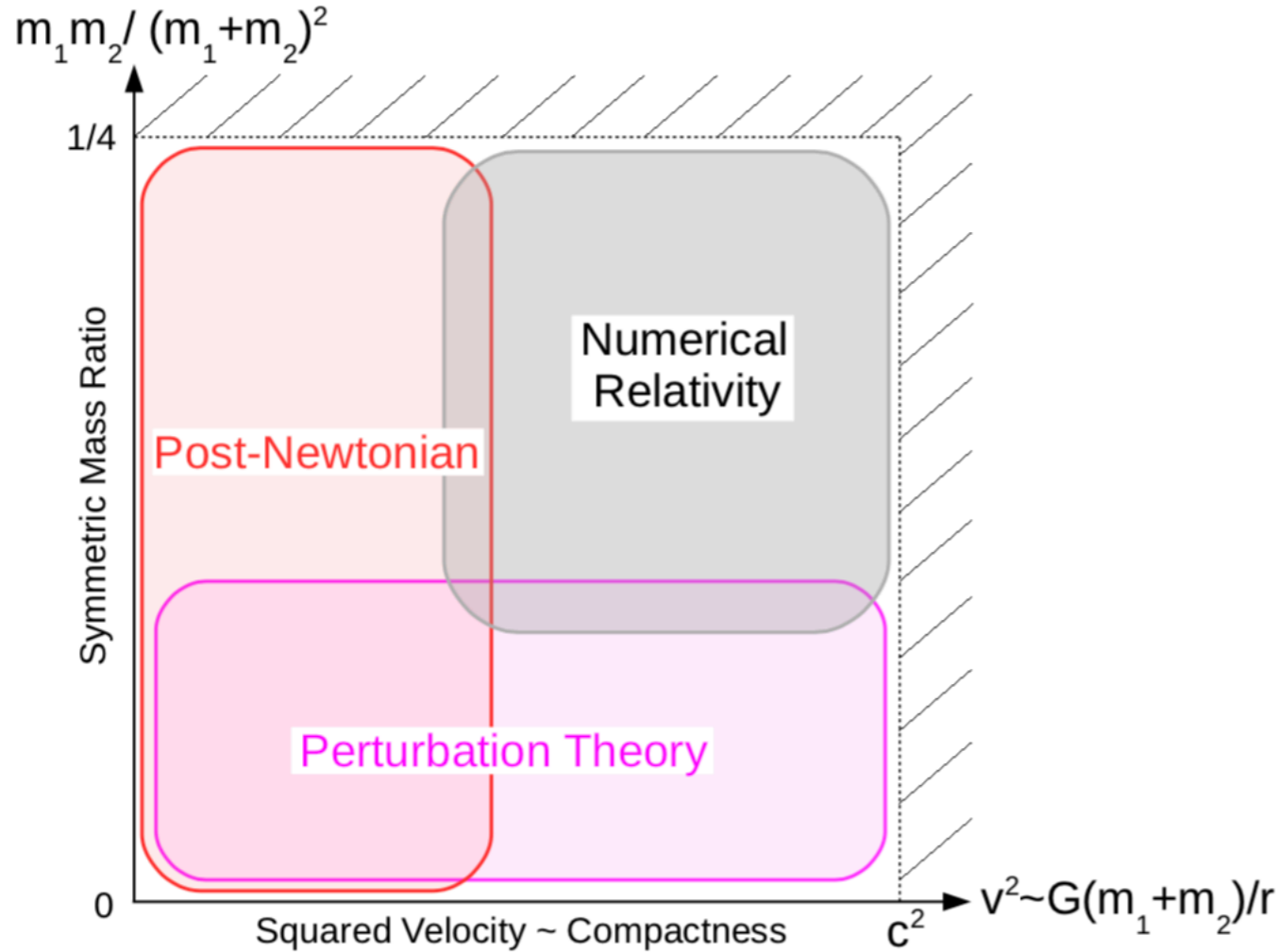


Part III : modeling the waveform

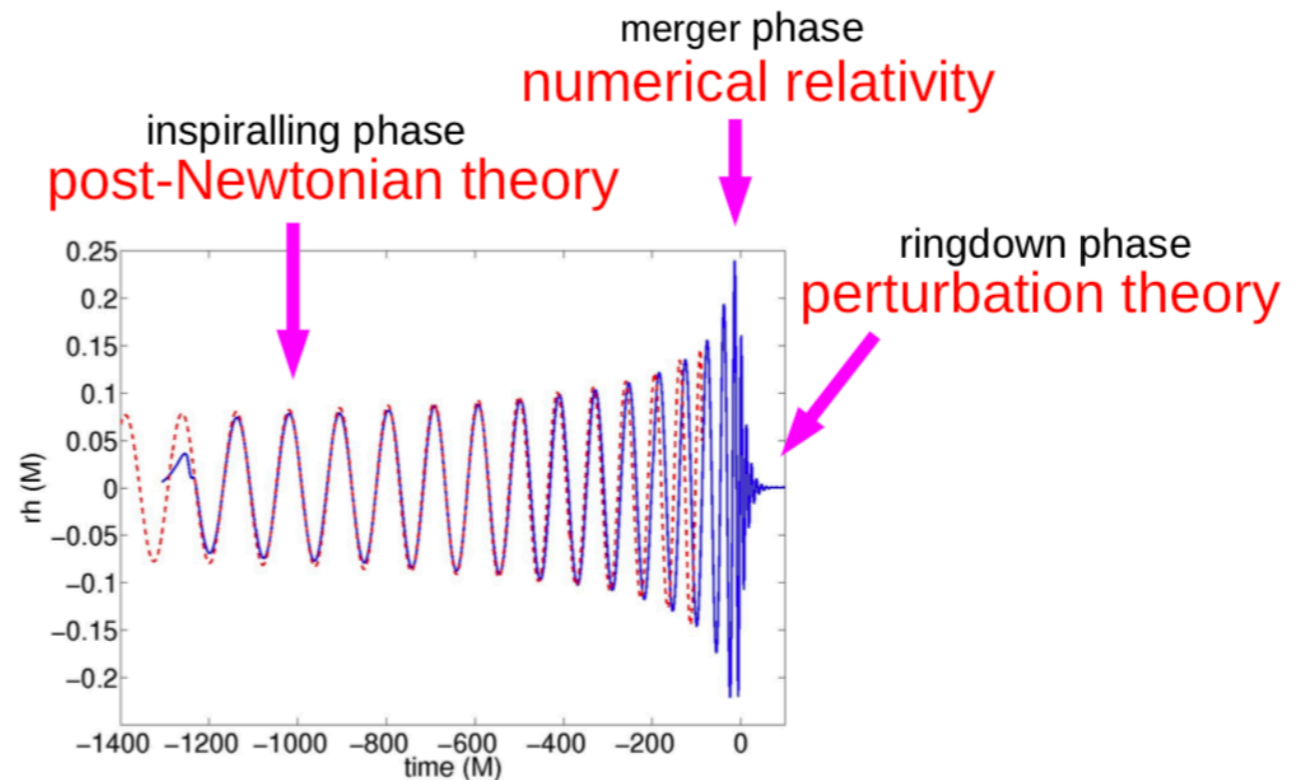
The need for a high-precision template



The different methods



(Credit : Luc Blanchet)



Post-Newtonian approximation

Perturbative solution of the EOM

$$G_{\mu\nu} = 8\pi T_{\mu\nu} \text{ and } g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}, \quad h_{\mu\nu} \sim \frac{GM}{r} \sim v^2 \ll 1$$

Matter is modeled by point particles (finite size effects arise only at $\mathcal{O}(v^8)$!):

$$S_m = -m_1 \int d\tau_1 - m_2 \int d\tau_2$$

Then plug back $g_{\mu\nu}$ in the action and obtain the two-body relativistic Lagrangian :

$$L = \frac{1}{2}m_1 v_1^2 + \frac{1}{2}m_2 v_2^2 + \frac{Gm_1 m_2}{r} + \text{relativistic corrections}$$

A (small) part of the 3PN energy

$$\begin{aligned}
 E = & \frac{m_1 v_1^2}{2} - \frac{Gm_1 m_2}{2r_{12}} \\
 & + \frac{1}{c^2} \left\{ \frac{G^2 m_1^2 m_2}{2r_{12}^2} + \frac{3m_1 v_1^4}{8} + \frac{Gm_1 m_2}{r_{12}} \left(-\frac{1}{4}(n_{12} v_1)(n_{12} v_2) + \frac{3}{2}v_1^2 - \frac{7}{4}(v_1 v_2) \right) \right\} \\
 & + \frac{1}{c^4} \left\{ -\frac{G^3 m_1^3 m_2}{2r_{12}^3} - \frac{19G^3 m_1^2 m_2^2}{8r_{12}^3} + \frac{5m_1 v_1^6}{16} \right. \\
 & \quad + \frac{Gm_1 m_2}{r_{12}} \left(\frac{3}{8}(n_{12} v_1)^3(n_{12} v_2) + \frac{3}{16}(n_{12} v_1)^2(n_{12} v_2)^2 - \frac{9}{8}(n_{12} v_1)(n_{12} v_2)v_1^2 \right. \\
 & \quad \quad - \frac{13}{8}(n_{12} v_2)^2 v_1^2 + \frac{21}{8}v_1^4 + \frac{13}{8}(n_{12} v_1)^2(v_1 v_2) + \frac{3}{4}(n_{12} v_1)(n_{12} v_2)(v_1 v_2) \\
 & \quad \quad \left. - \frac{55}{8}v_1^2(v_1 v_2) + \frac{17}{8}(v_1 v_2)^2 + \frac{31}{16}v_1^2 v_2^2 \right) \\
 & \quad \left. + \frac{G^2 m_1^2 m_2}{r_{12}^2} \left(\frac{29}{4}(n_{12} v_1)^2 - \frac{13}{4}(n_{12} v_1)(n_{12} v_2) + \frac{1}{2}(n_{12} v_2)^2 - \frac{3}{2}v_1^2 + \frac{7}{4}v_2^2 \right) \right\} \\
 & + \frac{1}{c^6} \left\{ \frac{35m_1 v_1^8}{128} \right. \\
 & \quad + \frac{Gm_1 m_2}{r_{12}} \left(-\frac{5}{16}(n_{12} v_1)^5(n_{12} v_2) - \frac{5}{16}(n_{12} v_1)^4(n_{12} v_2)^2 - \frac{5}{32}(n_{12} v_1)^3(n_{12} v_2)^3 \right. \\
 & \quad \quad + \frac{19}{16}(n_{12} v_1)^3(n_{12} v_2)v_1^2 + \frac{15}{16}(n_{12} v_1)^2(n_{12} v_2)^2 v_1^2 + \frac{3}{4}(n_{12} v_1)(n_{12} v_2)^3 v_1^2 \\
 & \quad \quad + \frac{19}{16}(n_{12} v_2)^4 v_1^2 - \frac{21}{16}(n_{12} v_1)(n_{12} v_2)v_1^4 - 2(n_{12} v_2)^2 v_1^4 \\
 & \quad \quad + \frac{55}{16}v_1^6 - \frac{19}{16}(n_{12} v_1)^4(v_1 v_2) - (n_{12} v_1)^3(n_{12} v_2)(v_1 v_2) \\
 & \quad \quad - \frac{15}{32}(n_{12} v_1)^2(n_{12} v_2)^2(v_1 v_2) + \frac{45}{16}(n_{12} v_1)^2 v_1^2(v_1 v_2) \\
 & \quad \quad \left. + \frac{5}{4}(n_{12} v_1)(n_{12} v_2)v_1^2(v_1 v_2) + \frac{11}{4}(n_{12} v_2)^2 v_1^2(v_1 v_2) - \frac{139}{16}v_1^4(v_1 v_2) \right) \\
 & \quad \left. + \frac{G^2 m_1^2 m_2}{r_{12}^2} \left(\frac{29}{4}(n_{12} v_1)^2 - \frac{13}{4}(n_{12} v_1)(n_{12} v_2) + \frac{1}{2}(n_{12} v_2)^2 - \frac{3}{2}v_1^2 + \frac{7}{4}v_2^2 \right) \right\}
 \end{aligned}$$

(Source : Luc Blanchet)

Dissipative dynamics

Similarly complicated calculations lead to the final result (for circular orbits) :

$$\phi = -\frac{x^{-5/2}}{32\nu} \left\{ 1 + \left(\frac{3715}{1008} + \frac{55}{12}\nu \right) x - 10\pi x^{3/2} + \left(\frac{15293365}{1016064} + \frac{27145}{1008}\nu + \frac{3085}{144}\nu^2 \right) x^2 \right. \\ + \left(\frac{38645}{1344} - \frac{65}{16}\nu \right) \pi x^{5/2} \ln\left(\frac{x}{x_0}\right) \\ + \left[\frac{12348611926451}{18776862720} - \frac{160}{3}\pi^2 - \frac{1712}{21}C - \frac{856}{21}\ln(16x) \right. \\ \left. + \left(-\frac{15737765635}{12192768} + \frac{2255}{48}\pi^2 \right) \nu + \frac{76055}{6912}\nu^2 - \frac{127825}{5184}\nu^3 \right] x^3 \\ \left. + \left(\frac{77096675}{2032128} + \frac{378515}{12096}\nu - \frac{74045}{6048}\nu^2 \right) \pi x^{7/2} + \mathcal{O}\left(\frac{1}{c^8}\right) \right\},$$

$$x = (GM\omega)^{2/3} \sim v^2$$

$$\nu = \frac{m_1 m_2}{(m_1 + m_2)^2}$$

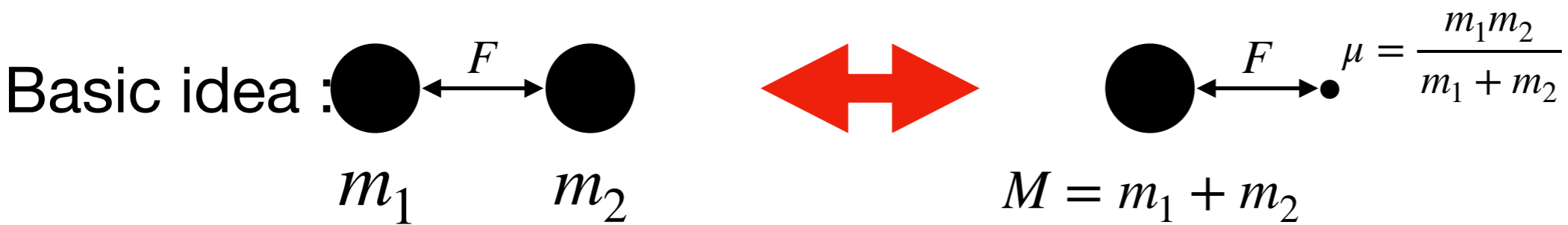
	$2 \times 1.4 M_\odot$	$10 M_\odot + 1.4 M_\odot$	$2 \times 10 M_\odot$
Newtonian order	16031	3576	602
1PN	441	213	59
1.5PN (dominant tail)	-211	-181	-51
2PN	9.9	9.8	4.1
2.5PN	-11.7	-20.0	-7.1
3PN	2.6	2.3	2.2
3.5PN	-0.9	-1.8	-0.8

(source : Luc Blanchet)

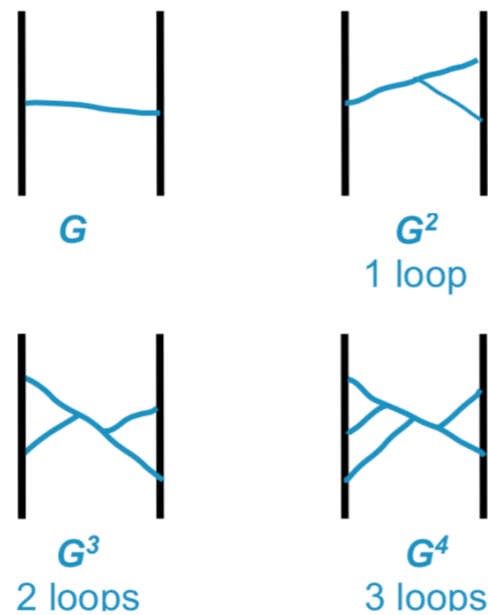
Note the slow convergence !

Effective One-Body (Buonanno-Damour 98)

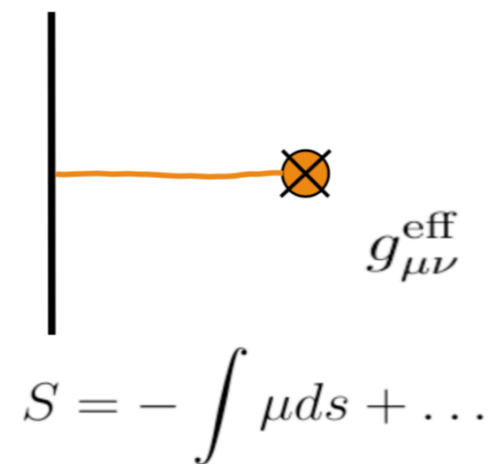
We want to resum as much as possible the (badly convergent) PN expansion



Real dynamics



Effective dynamics



$$H = H_0 + \left(GH_1 + \frac{G^2}{c^2} H_2 + \frac{G^3}{c^4} H_3 + \frac{G^4}{c^6} H_4 \right) \left(1 + \frac{1}{c^2} + \dots \right)$$

$$A(R) = 1 - \frac{2GM}{c^2 R} + 2\nu \left(\frac{GM}{c^2 R} \right)^3 + \dots,$$

$$B(R) = 1 + \frac{2GM}{c^2 R} + (4 - 6\nu) \left(\frac{GM}{c^2 R} \right)^2 + \dots$$

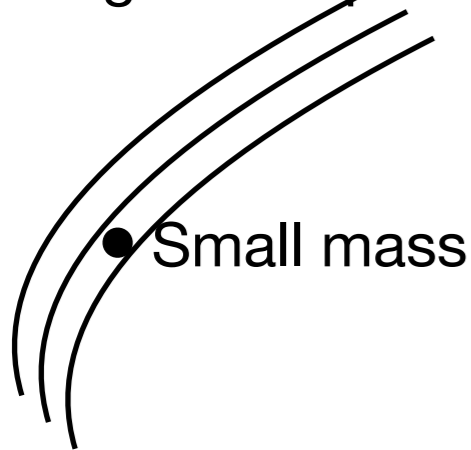
Effective metric

$$ds_{\text{eff}}^2 = -A(r)dt^2 + B(r)dr^2 + r^2 (d\theta^2 + \sin^2 \theta d\varphi^2)$$

Gravitational self-force

EMRI : find the waveform as an expansion in $\epsilon = m/M$. This is far from being achieved.

Background spacetime



The tensorial form of the equations of motion is

$$\frac{Du^\mu}{d\tau} = -\frac{1}{2}(g^{\mu\nu} + u^\mu u^\nu)(2h_{\nu\lambda\rho}^{\text{tail}} - h_{\lambda\rho\nu}^{\text{tail}})u^\lambda u^\rho$$

where

Deviation from geodesics

$$h_{\mu\nu\lambda}^{\text{tail}} = 4m \int_{-\infty}^{\tau-\epsilon} \nabla_\lambda \left(G_{\mu\nu\mu'\nu'} - \frac{1}{2}g_{\mu\nu}G^\rho{}_{\rho\mu'\nu'} \right) (z(\tau), z(\tau')) u^{\mu'} u^{\nu'} d\tau'$$

(MiSaTaQuWa equations)

Mathematically challenging to go to higher orders, but $\mathcal{O}(\epsilon^3)$ is needed !

Conclusions

- We will now be able for the first time to observe the universe with another channel than photons !
- We begin to observe the close universe, and we will observe it to cosmological distances in few decades
- Modeling of the signal is analytically and numerically challenging



Il n'y a pas de problèmes résolus, seulement des problèmes plus ou moins résolus